

## Sensing extremely limited H<sub>2</sub> contents by Pd nanogap connected to an amorphous InGaZnO thin-film transistor†

Cite this: *Nanoscale*, 2013, 5, 8915

Received 13th April 2013  
Accepted 15th July 2013

DOI: 10.1039/c3nr01847d

www.rsc.org/nanoscale

Young Tack Lee,<sup>‡a</sup> Hwaebong Jung,<sup>‡b</sup> Seung Hee Nam,<sup>ac</sup> Pyo Jin Jeon,<sup>a</sup>  
Jin Sung Kim,<sup>a</sup> Byungjin Jang,<sup>b</sup> Wooyoung Lee<sup>\*b</sup> and Seongil Im<sup>\*a</sup>

A palladium (Pd) nanogap-based thin-film has been connected to an electrically stable amorphous InGaZnO thin-film transistor, to form a hydrogen sensor demonstrating a dramatic sensing capability. As a result of the Pd connection to the transistor source, our sensor circuit greatly enhances the hydrogen-induced signal and sensing speed in the sense of output voltage, clearly resolving a minimum hydrogen content of 0.05%. When the nanogap-based Pd thin-film was connected to the transistor gate, an extremely limited hydrogen content of even less than 0.05% was visibly detected through gate voltage shifts. Our results exhibit the most promising and practical ways to sense extremely limited hydrogen contents, originating from two methods: transistor-to-Pd nanogap resistor and transistor-to-Pd nanogap capacitor coupling.

Hydrogen sensing has continually drawn great attention since it was regarded important in industrial safety, which claims that 3–4% of ambient H<sub>2</sub> is the margin to avoid explosion danger.<sup>1–9</sup> So, monitoring the ambient H<sub>2</sub> and detecting an even extremely small content of H<sub>2</sub> gas is not only important but also always requested, particularly with practical and simple sensing tools.<sup>6,10–23</sup> Very recently, we reported one of the most innovative H<sub>2</sub> sensors, using typical Pd thin-films but containing nanogaps which were created by mechanically stretching an elastomeric substrate.<sup>20–23</sup> The sensitivity of our nanogap-based sensor appeared very good, so that less than ~0.1% of H<sub>2</sub> gas molecules in a N<sub>2</sub> atmosphere might be detected by reading signal current. However, reading the detection current at such low H<sub>2</sub> content was not easy at all; in fact even 0.2% hydrogen gas was barely noted by current sensors. More scientific techniques to

improve the detection capability is thus necessary. In the present work, we dramatically improved the H<sub>2</sub> sensing capability of the nanogap-based Pd sensor by connecting it to an electrically stable amorphous InGaZnO thin-film transistor (a-IGZO TFT) in two different ways: Pd connection to the TFT source and to the gate. The IGZO TFT was chosen, since it was stable enough to bear the gate bias stress during H<sub>2</sub> detection; it would eventually be integrated with our Pd sensor. As a result of the Pd connection to the TFT source, our sensor circuit greatly enhances the H<sub>2</sub>-induced signal by three orders of magnitude in the sense of output voltage, clearly resolving a minimum hydrogen content of 0.05%. When the nanogap-based Pd sensor was connected to the TFT gate, an even lower H<sub>2</sub> content of less than 0.05% was visibly detected, but more interestingly through this study of TFT gate oxide capacitor-to-Pd nanogap capacitor coupling, the size of the nanogap distance was approximately estimated as an important result.

Polydimethylsiloxane (PDMS) elastomer was used as the substrate of the hydrogen gas sensor, to create nanocracks (or nanogap stripes) in the Pd thin-film (TF). The base resin (Sylgard 184, Dow Corning) was mixed with a curing agent at a volume ratio of 10 : 1 and was kept in the vacuum chamber for 10 min to evacuate any air bubbles. After then, the PDMS mixture was cured for 30 min on a hotplate at 423 K. In order to fabricate the hydrogen gas sensor, a 10 nm-thin Pd film was deposited on the PDMS substrate by using a DC magnetron sputtering system. After Pd deposition, the Pd–PDMS substrate was mounted onto a stretching machine and nanocracks were created under a tensile stress.<sup>20–23</sup> (ESI, Fig. S1†) Fig. 1a shows a schematic of the as-cracked Pd TF device (area dimension: 10 mm [L] × 20 mm [W]), where many nanogap stripes are shown and one of them is illustrated in Fig. 1b with its schematic cross section. The single nanocrack appears deformed but yet maintains an internal connection (the connection was confirmed by an electric conduction measurement). However, if the initially cracked region comes into contact with H<sub>2</sub> gas molecules (*i.e.* is exposed to H<sub>2</sub> to form a PdH<sub>x</sub> compound; Fig. 1c) and the molecules are evacuated, the crack is left

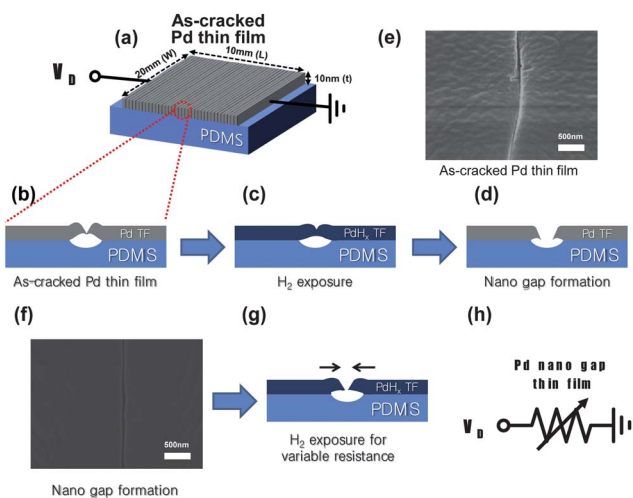
<sup>a</sup>Institute of Physics and Applied Physics, Yonsei University, Seoul 120-749, Korea. E-mail: semicon@yonsei.ac.kr; Fax: +82-2-392-1592; Tel: +82-2-2123-2842

<sup>b</sup>Department of Materials Science and Engineering, Yonsei University, Seoul 120-749, Korea. E-mail: Wooyoung@yonsei.ac.kr; Fax: +82-2-312-5275; Tel: +82-2123-2834

<sup>c</sup>R&D Center, LG Display, Paju 413-811, Korea

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c3nr01847d

‡ These authors contributed equally to this work.

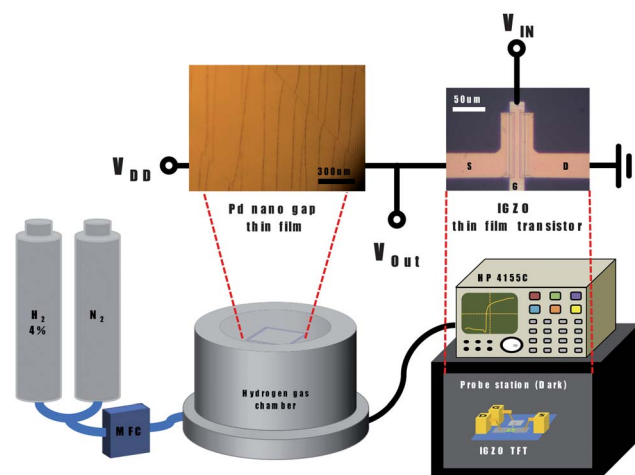


**Fig. 1** Schematics of (a) an as-cracked Pd TF device, (b) an as-cracked nanogap stripe in cross sectional view, (c) a PdH<sub>x</sub> compound-containing stripe as exposed to H<sub>2</sub>, and (d) the stripe with an eternally open nanogap. SEM images of (e) and (f) show a stripe from the as-cracked and eternal nanogap-containing Pd TF, respectively. (g) A schematic cross section of a nanogap stripe in Pd TF with variable electrical resistances which depend on H<sub>2</sub> exposure, and (h) its equivalent circuit.

eternally open with a nanogap as shown in Fig. 1d. Fig. 1e and f are the respective scanning electron microscopy (SEM) images of as-cracked and eternal nanogap-containing Pd TF. According to the SEM images, the as-cracked Pd still shows a physical contact between the two Pd regions around the crack stripe while the other stripe from nanogap-containing Pd shows ~50 nm separation (gap distance), although the distance is not uniform. Based on the structure of Fig. 1d and f, the nanogap Pd TF now operates as a H<sub>2</sub> sensor, since ON/OFF (contact/no contact) switching takes place in the Pd TF with variable resistances according to the degree of H<sub>2</sub> molecule adsorption on the Pd surface or H<sub>2</sub> reaction with Pd. The H<sub>2</sub>-adsorbed Pd TF changes its phase to PdH<sub>x</sub> with an expanded volume which leads to a physical contact in the nanogap (see the illustration in ESI, Fig. S2†).<sup>24–27</sup> As a result, the variable electrical conductance of nanogap Pd TF may indicate H<sub>2</sub> content (see Fig. 1g and h for the illustration and circuit of the variable resistance).

The structure of a-IGZO TFTs is an inverted-stagger type with a width-to-length (*W/L*) ratio of 100 : 10 μm using a bottom gate (ESI, Fig. S3†). A 300 nm-thick SiO<sub>2</sub> gate insulator layer was deposited on the patterned Cu–MoTi gate electrode by a plasma-enhanced chemical vapor deposition (PECVD) system. Then, a 60 nm-thick active channel layer (a-IGZO) was deposited by a DC magnetron sputtering system. The formation of the SiO<sub>2</sub> etch stopper (75 nm) and Mo source/drain electrodes were sequentially performed, followed by the formation of a 300 nm-thick PECVD SiO<sub>2</sub> passivation layer. All the patterning processes were carried out by photolithography that involved wet chemical etching and PECVD processes. The device annealing at 300 °C in ambient air was performed as a final process.

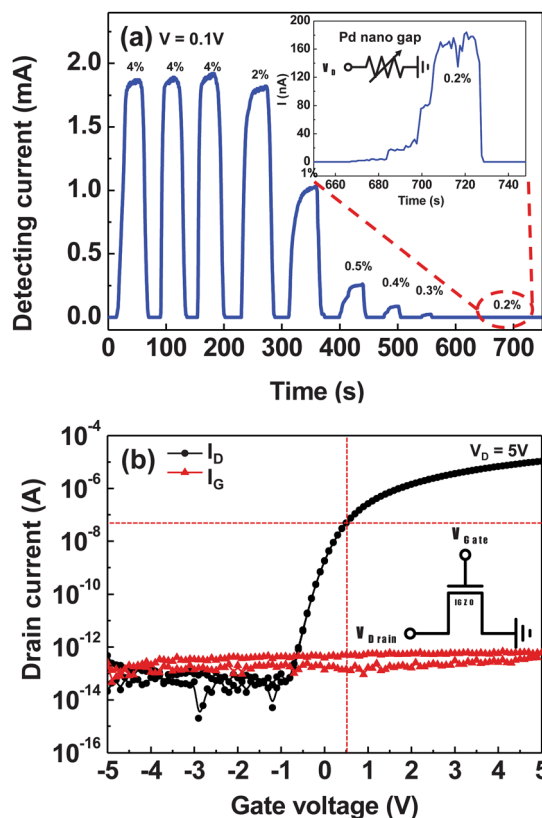
Fig. 2 shows a schematic diagram of our H<sub>2</sub> sensor system, which consists of a small gas chamber with a volume of ~250 ml and a mass flow controller (MFC) for pure 100% N<sub>2</sub> (for



**Fig. 2** Schematic illustration of our H<sub>2</sub> sensor system. The nanogap Pd TF device was mounted in the gas chamber and connected to an electrically stable a-IGZO TFT device of a probe station in the dark. Optical microscope images for a nanogap-stripe Pd TF resistor and TFT are presented in the inset type circuit.

purging) and 96% N<sub>2</sub> + 4% H<sub>2</sub> mixture gas (for sensing). The pressure in the chamber was maintained at 1 atm at room temperature. The nanogap Pd TF device was mounted in the gas chamber and connected to an electrically stable a-IGZO TFT device (to the source electrode) of a probe station in the dark. Supply and input voltages (*V*<sub>DD</sub> and *V*<sub>IN</sub>) were applied by a semiconductor parameter analyzer (Model HP 4155C, Agilent Technologies). Output voltage (*V*<sub>OUT</sub>) according to H<sub>2</sub> sensing was also measured by the semiconductor analyzer. The circuit illustration in Fig. 2 contains the top view images of our nanogap Pd TF device and a-IGZO TFT, which were obtained from an optical microscope. As shown, our Pd TF contains many nanogap stripes, of which the number was approximated to be 60–70. Our measurement setup is analogous to a logic inverter circuit for digital output, except that ours has a variable resistor composed of a H<sub>2</sub>-sensing nanogap Pd TF instead of a constant load resistor; we could thus expect that small H<sub>2</sub>-induced analogue current would be properly amplified in voltage signals.

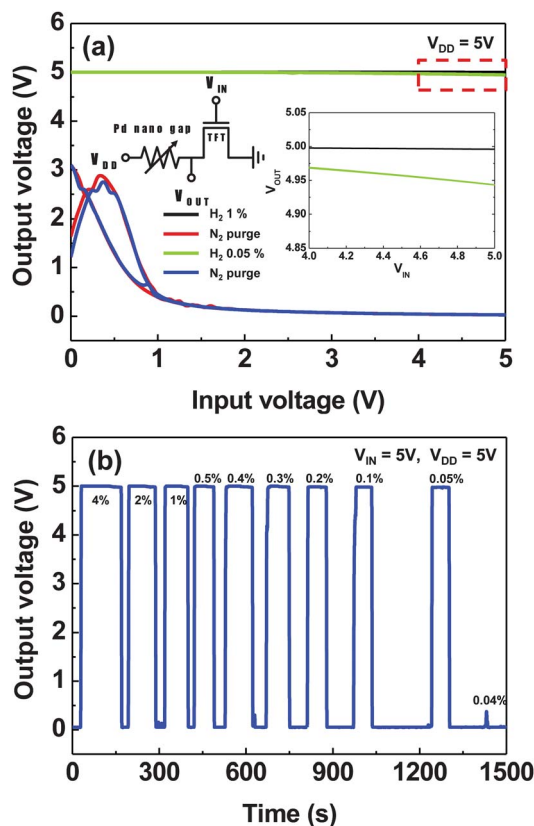
Fig. 3a is a detecting current vs. time plot which is obtained from our nanogap Pd TF (under 0.1 V) by H<sub>2</sub> gas-controlled ON/OFF switching. Various H<sub>2</sub> ambiences of 4%, 2%, 1%, 0.5%, 0.4%, 0.3%, and 0.2% (as mixed with pure N<sub>2</sub>) were exposed to the Pd sensor. While H<sub>2</sub> concentrations higher than 0.3% were quite nicely detected at the current range from ~2 mA to few hundred μA, it was not easy to note a lower H<sub>2</sub> content from the current signal. In particular, 0.2% H<sub>2</sub> was barely sensed and noted only by magnifying the sensing signal (~200 nA) as shown in the inset of Fig. 3a. Worse than that, sensing (and recovery) time gets longer with smaller H<sub>2</sub> content, so that detecting 0.3% H<sub>2</sub> appeared to take more than 20 s. Hence, we employed an electrically stable a-IGZO TFT to amplify the sensing signal and simultaneously to shorten the sensing time, converting the output signal from current to voltage. Fig. 3b is the drain current–gate voltage (*I*<sub>D</sub>–*V*<sub>G</sub>) transfer characteristics of our oxide TFT along with gate leakage current (*I*<sub>G</sub>) curves, as



**Fig. 3** (a) Detecting current vs. time plot achieved by H<sub>2</sub> gas-controlled ON/OFF switching in nanogap Pd TF. Various H<sub>2</sub> contents of 4%, 2%, 1%, 0.5%, 0.4%, 0.3%, and 0.2% were exposed to the Pd sensor and purged by pure N<sub>2</sub>. Inset shows the current for 0.2% H<sub>2</sub> in zoomed scale. (b) Drain current–gate voltage ( $I_D$ – $V_G$ ) transfer characteristics of our a-IGZO TFT device.

obtained under a drain bias of  $V_D = 5$  V. As shown in the TFT transfer curves of Fig. 3b, our TFT was very stable with 0.1 pA  $I_G$  leakage and without any gate hysteresis, turning on above 0 V. (ON/OFF ratio was more than  $\sim 108$  and the saturation field-effect mobility was  $\sim 10$  cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>). Since the transistor  $I_D$  current covers a broad range from 20  $\mu$ A to 0.1 pA while our Pd sensor has another current range from 2 mA to a few hundred nA, some reasonable  $V_{OUT}$  signals are expected through an overlapped current region when these two devices are connected in series.

The series connection scheme has already been introduced in Fig. 2 and is now shown in the inset circuit diagram of Fig. 4a, where  $V_{OUT}$  vs.  $V_{IN}$  plots are displayed as the main measurement results. For H<sub>2</sub> detection experiments, H<sub>2</sub> gas of 1%, 0.05%, and 0% (purged with 100% N<sub>2</sub> gas) was infused into the hydrogen sense chamber of Fig. 2 in the following sequence: initial 1% H<sub>2</sub> infusion, N<sub>2</sub> purge, 0.05% H<sub>2</sub> infusion, and final N<sub>2</sub> purge, so that each  $V_{OUT}$  curve could be obtained under a  $V_{DD} = V_{IN} = 5$  V as shown in Fig. 4a. According to the plots, the  $V_{OUT}$  for 0.05% H<sub>2</sub> detection appears almost the same as that for 1% H<sub>2</sub> sensing, indicating almost 5 V; this result is confirmed by magnifying some part of the plot (see the inset plot for a range of  $V_{IN} = 4$ –5 V). Without H<sub>2</sub> (by N<sub>2</sub> purge),  $V_{OUT}$  drops to 0 V at a high  $V_{IN}$ , however the drop stops near  $\sim 2.5$  V at a low  $V_{IN}$  of 0.5 V, which



**Fig. 4** (a)  $V_{OUT}$  vs.  $V_{IN}$  plots of our inverter type H<sub>2</sub> sensor tested with 1%, 0.05%, and 0% (pure N<sub>2</sub>) of H<sub>2</sub> contents. The inset shows the  $V_{OUT}$  curves for 1 and 0.05% H<sub>2</sub> in zoomed scale. (b) Dynamic H<sub>2</sub> detection plot in time domain for various H<sub>2</sub> contents from 4% to 0.05%.

might be too low to open the transistor channel properly. Based on these  $V_{OUT}$ – $V_{IN}$  plots, we implemented a time domain H<sub>2</sub> detection experiment (for  $V_{OUT}$  vs. time plot), dynamically varying the H<sub>2</sub> concentration from 4% to 0.04% under a fixed voltage condition of  $V_{DD} = V_{IN} = 5$  V. Fig. 4b displays the dynamic H<sub>2</sub> detection plot in time domain, where the detection voltage  $V_{OUT}$  appears to be  $\sim 5$  V for all the cases from 4% to 0.05% H<sub>2</sub>, as predicted by the results from Fig. 4a. (The results in Fig. 4b were reproducible and a magnified time domain  $V_{OUT}$  plot is seen in the ESI, Fig. S4.†) The detection voltage abruptly decreases to 0.4 V from  $\sim 5$  V when the H<sub>2</sub> content further decreases to 0.04%, although this specific result was actually irreproducible unlike the case of 0.05% and higher H<sub>2</sub> percentages. We thus recognize that the series connection setup for H<sub>2</sub> sensing still has its own detection limit at 0.05%. In spite of such a limitation, however, it is acknowledgeable that applying our logic inverter type circuit to the analogue nanogap Pd TF remarkably enhances the visibility of the low % H<sub>2</sub> signal by a few orders of magnitude and simultaneously shortens the sensing (and recovery) time to 1–2 s. In fact, the electrical conductance or conducting path in nanogap Pd film still exists even with the low 0.05% hydrogen molecules, but is too small to measure/or estimate and thus becomes visible only with the connection to a transistor. Apparent electrical disconnection in the nanogap Pd film is eventually observed with an extremely

low % of H<sub>2</sub> molecules (here 0.04% is the onset point of the open circuit).

Complete gap opening might take place in single or a few nanogap stripes at a certain point of dilute ambient H<sub>2</sub>, leading to an abrupt disconnection. In this event, our nanogap-containing Pd TF has then a charge capacitance due to the gap distance. The gap distance would be so dependent on the H<sub>2</sub> content that we may estimate an extremely low % of H<sub>2</sub> by introducing the capacitance of the nanogap Pd TF to the gate of the same a-IGZO TFT with a dielectric oxide (300 nm thick SiO<sub>2</sub>). Such a capacitor-induced circuit is shown in the inset of Fig. 5a and the transfer curves have been accordingly obtained from the circuit with the oxide TFT and Pd TF variable capacitor under two ambient conditions of 0.04% and 0% H<sub>2</sub> (by N<sub>2</sub> purge). The initial transfer curve was obtained as a reference from the TFT alone. Connected in series to the Pd TF, the total gate capacitance ( $C_{\text{total}}$ ) of the inset circuit should become smaller than the dielectric oxide capacitance  $C_{\text{ox}}$  because a nanogap-induced capacitance ( $C_{\text{gap}}$ ) should be also considered as shown below in eqn (1).

$$C_{\text{total}} = \frac{1}{\left(\frac{1}{C_{\text{gap}}} + \frac{1}{C_{\text{ox}}}\right)} \quad (1)$$

The smaller  $C_{\text{total}}$  leads to a slightly smaller  $I_{\text{D}}$  current than that by  $C_{\text{ox}}$  alone at an identical on-state  $V_{\text{G}}$  (e.g. 3 V), as shown in Fig. 5a. Simultaneously, at the same off-state  $I_{\text{D}}$  (e.g.  $\sim 0.3$  pA) the smaller  $C_{\text{total}}$  leads to a more negative  $V_{\text{G}}$ , that is, the transfer curves of the TFT with the Pd TF capacitor shows more negatively shifted  $V_{\text{G}}$  with a lower H<sub>2</sub> content, since the  $C_{\text{gap}}$  makes an additional charging and voltage drop ( $V_{\text{gap}} = V_{\text{G}} - V_{\text{Gi}}$ ) prior to  $C_{\text{ox}}$ -induced gate charging. The  $C_{\text{gap}}$  can be expressed as  $A\varepsilon_0/d$ , where  $A$  is the thickness of the wall area of a nanogap stripe (10 nm  $\times$  20 nm),  $d$  is the gap distance, and  $\varepsilon_0$  is the

dielectric constant in air (or N<sub>2</sub>). The  $V_{\text{gap}}$  is measured from the transfer curves where an off-state voltage ( $V_{\text{Gi}}$ ) of the initial TFT is indicated by a dashed line of Fig. 5a. These  $V_{\text{G}}$  shift results may provide a key solution for sensing even an extremely low H<sub>2</sub>% and simultaneously for the estimation of the H<sub>2</sub>%-dependent average gap distance as well. Using the  $V_{\text{gap}}$  and  $C_{\text{gap}}$  dependence on H<sub>2</sub>%, we estimated the average gap distances in 0% and 0.04% H<sub>2</sub> ambiances to be  $\sim 30$  nm and  $\sim 9$  nm, respectively. Since the circuit is connected to the gate through the Pd TF resistor anyway if the H<sub>2</sub> content in the chamber atmosphere is over 0.05%, the transfer curves under such conditions become almost identical to the initial transfer curve obtained without the Pd TF connection; in fact Fig. 5b displays the two curves quite overlapped on the initial curve. An equivalent circuit is shown in the inset of Fig. 5b, where the nanogap Pd TF is expressed as a variable resistor. Fig. 5c and d are the respective illustrations of Pd TF capacitor and resistor models with diluted and concentrated H<sub>2</sub> molecules in a Pd lattice.

Merging the TFT gate oxide and nanogap Pd TF capacitors, the total gate capacitance ( $C_{\text{total}}$ ) becomes smaller than the dielectric oxide capacitance (here  $C_{\text{ox}}$  for 300 nm thick SiO<sub>2</sub>). We already addressed this point in eqn (1), since a nanogap-induced capacitance ( $C_{\text{gap}}$ ) should be considered in addition to  $C_{\text{ox}}$ . At the same off-state  $I_{\text{D}}$  (e.g.  $\sim 0.3$  pA) an open circuit-based transfer curve shows more negatively shifted  $V_{\text{G}}$ , since the  $C_{\text{gap}}$  makes an additional charging and voltage drop ( $V_{\text{gap}} = V_{\text{G}} - V_{\text{Gi}}$ ) prior to  $C_{\text{ox}}$ -induced gate charging. A schematic circuit of this situation is illustrated in Fig. 6. The following eqn (2)–(4) are now useful for the estimation of the nanogap distance. Due to a series capacitance theory in general physics,

$$Q_{\text{gap}} = Q_{\text{ox}} = Q_{\text{s}}, \quad (2)$$

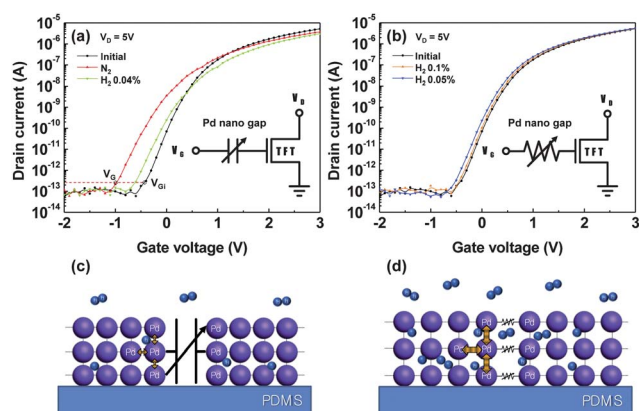
and

$$C_{\text{gap}}V_{\text{gap}} = C_{\text{ox}}V_{\text{ox}} = 2C_{\text{s}}\Psi_{\text{s}} \quad (3)$$

then,

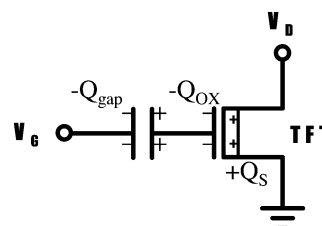
$$V_{\text{Gi}} - V_{\text{FB}} = V_{\text{ox}} + \Psi_{\text{s}} + V_{\text{c}}, \quad (4)$$

where  $Q_{\text{gap}} (= C_{\text{gap}}V_{\text{gap}})$ ,  $Q_{\text{ox}} (= C_{\text{ox}}V_{\text{ox}})$ , and  $Q_{\text{s}} (= 2C_{\text{s}}\Psi_{\text{s}})$  are total charges in the nanogap, oxide–gate interface, and depleted a-IGZO film, respectively. Eqn (4) is from the TFT without the Pd TF capacitor.  $V_{\text{Gi}}$  is an initial off-state gate voltage of the TFT



**Fig. 5** (a)  $I_{\text{D}}-V_{\text{G}}$  Transfer curves obtained from the initial a-IGZO TFT (black line) and from the a-IGZO TFT linked to the Pd TF variable capacitor under two ambient conditions: 0% H<sub>2</sub> (pure N<sub>2</sub>; red line) and 0.04% H<sub>2</sub> (green line). Inset is the corresponding equivalent circuit with the capacitor. (b)  $I_{\text{D}}-V_{\text{G}}$  transfer curves obtained from the a-IGZO TFT device (black line) and from the a-IGZO TFT linked to the Pd TF variable resistor in another two ambiances: 0.1% (orange line) and 0.05% H<sub>2</sub> (blue line). Inset is the corresponding circuit with the resistor. Illustrations (c) and (d) are the respective models of Pd TF capacitor and resistor with diluted and concentrated H<sub>2</sub> molecules in a Pd lattice.

### Nanogap Pd TF + a-IGZO TFT circuit



**Fig. 6** A circuit to express the series capacitance theory, with eqn (2) and (3) in a merged system of the TFT gate oxide and nanogap Pd TF capacitor.



without the Pd TF capacitor,  $V_{\text{FB}}$  is the flat band voltage,  $\Psi_s$  is a voltage for semiconductor band bending for an off-state (semiconductor charge depletion),  $V_{\text{ox}}$  is the band bending of the dielectric  $\text{SiO}_2$ ,  $V_c$  is an average channel voltage ( $\langle V_D = 5 \text{ V} \rangle$ ), and  $C_s$  is the capacitance of a charge-depleted 60 nm-thin a-IGZO semiconductor. (Note that  $Q_{\text{gap}} (= C_{\text{gap}}V_{\text{gap}})$  and  $Q_{\text{ox}} (= C_{\text{ox}}V_{\text{ox}})$  are sheet charges but  $Q_s (= 2C_s\Psi_s)$  is the charge in a thin film; that is why the factor of 2 is added.)<sup>28</sup> Whether the Pd TF capacitor is connected or not, the same band bending terms ( $V_{\text{ox}}$  and  $\Psi_s$ ) should be maintained in both cases for a transistor (with and without the Pd TF capacitor) to keep an identical off-state current  $I_D$  ( $\sim 3 \text{ pA}$ ). Now, from the capacitance–voltage ( $C$ – $V$ ) curve of Fig. 7, we may extract a part of eqn (4),  $V_{\text{ox}} + \Psi_s$  since in  $C$ – $V$  measurements  $V_{\text{FB}}$  is approximated by the cross point of two tangent lines while  $V_c$  is zero so that  $V_{\text{Gi}}$  becomes  $V_{\text{Go}}$ ;  $V_{\text{ox}} + \Psi_s$  is  $\sim -1.0 \text{ V}$  as estimated from  $V_{\text{Go}} - V_{\text{FB}}$  in the  $C$ – $V$  curve. The minimum  $C_{\text{total}}$  in the  $C$ – $V$  curve is observed to be  $\sim 0.025 \text{ pF}$  ( $= 1/(1/C_s + 1/C_{\text{ox}})$ , eqn (1)) in a depletion state and  $C_{\text{ox}}$  is measured to be  $\sim 0.131 \text{ pF}$  in an accumulation state of  $V_G$  (the inset band diagrams respectively indicate the depletion and accumulation states). As a result from eqn (1),  $C_s$  is calculated to be  $\sim 0.031 \text{ pF}$ , since we measure the  $C_{\text{ox}}$  from the  $C$ – $V$  curve. Then we use the relation  $C_{\text{ox}}V_{\text{ox}} = 2C_s\Psi_s$  along with another relation  $V_{\text{ox}} + \Psi_s = \sim -1.0 \text{ V}$ , to work out  $V_{\text{ox}}$  and  $\Psi_s$  ( $-0.32 \text{ V}$  and  $-0.68 \text{ V}$ , respectively). Since our Pd TF stripe area ( $A$ ) is  $10 \text{ nm} \times 20 \text{ nm}$  (for  $t \times W$ , see Fig. 1a) and it can be applied to eqn (3), we can finally estimate the gap distance using another equation below,

$$C_{\text{gap}} = \frac{\epsilon_0 A}{Nd} = \frac{C_{\text{ox}}V_{\text{ox}}}{V_{\text{gap}}}. \quad (5)$$

Now,

$$Nd = \frac{\epsilon_0 A V_{\text{gap}}}{C_{\text{ox}} V_{\text{ox}}}, \quad (6)$$

where  $N$  is the number of nanogap stripes playing as capacitors in Pd TF,  $d$  is the gap distance of a single capacitor,  $V_{\text{gap}}$  is  $0.7 \text{ V}$  for  $\text{N}_2$  purge conditions and  $0.22 \text{ V}$  for  $0.04\% \text{ H}_2$  according to Fig. 5a. Average  $Nd$  values are calculated to be  $\sim 30 \text{ nm}$  and

$\sim 9 \text{ nm}$  for the above two respective cases ( $0\%$  and  $0.04\% \text{ H}_2$ ).  $N$  might be 1 or more, but here it is very likely that  $N$  is 1 for the  $0.04\%$  case, since we have observed an abrupt electrical disconnection at  $0.04\% \text{ H}_2$  with only a little dilution of  $\text{H}_2$  from  $0.05$  to  $0.04\%$ . The maximum gap distance would be  $\sim 30 \text{ nm}$  as in the case of  $0\% \text{ H}_2$ . Our calculation of the gap distances may not provide an exact number value, however, the calculated value is at least comparable to the scale observed in the SEM results of Fig. 1f.

## Conclusions

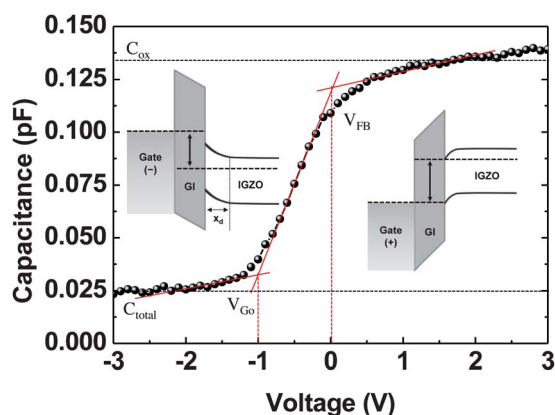
In summary, we have fabricated an ultrasensitive  $\text{H}_2$  sensor system composed of an electrically stable a-IGZO TFT and a nanogap-containing Pd TF, which acts as an analogue variable resistor depending on the  $\text{H}_2$  content. By connecting the nanogap Pd TF to the TFT source, our inverter type sensor circuit greatly enhances the  $\text{H}_2$ -induced signal visibility and sensing speed in a sense of output voltage, clearly resolving a minimum hydrogen content of  $0.05\%$ . When the nanogap-based Pd TF was connected to the TFT gate, an even lower  $\text{H}_2$  content of less than  $0.05\%$  was visibly detected, since the Pd TF acts as a capacitor with  $\text{H}_2$ -dependent variable capacitance. We thus conclude that our system of oxide TFT-coupled Pd TF with nanogaps is very promising for sensing an extremely low  $\text{H}_2$  content in the ambient atmosphere.

## Acknowledgements

Authors acknowledge the financial support from NRF grant (NRL program: Grant no. 2009-0079462) and Priority Research Centers Program (2009-0093823) through the National Research Foundation of Korea (NRF). Young Tack Lee would like to thank for the LGenius Members Fellowship from LGD company.

## Notes and references

- 1 A. Katsuki and K. Fukui, *Sens. Actuators, B*, 1998, **52**, 30.
- 2 K. I. Lundström, M. S. Shivaraman and C. M. Svensson, *J. Appl. Phys.*, 1975, **46**, 3876.
- 3 B. S. Kang, R. Mehandru, S. Kim, F. Ren, R. C. Fitch, J. K. Gillespie, N. Moser, G. Jessen, T. Jenkins, R. Dettmer, D. Via, A. Crespo, B. P. Gila, C. R. Abernathy and S. J. Pearton, *Appl. Phys. Lett.*, 2004, **84**, 4635.
- 4 M. Yun, N. V. Myung, R. P. Vasquez, C. Lee, E. J. Menke and R. M. Penner, *Nano Lett.*, 2004, **4**, 419.
- 5 Y. Im, C. Lee, R. P. Vasquez, M. A. Bangar, N. V. Myung, E. J. Menke, R. M. Penner and M. Yun, *Small*, 2006, **2**, 356.
- 6 F. Favier, E. C. Walter, M. P. Zach, T. Benter and R. M. Penner, *Science*, 2001, **293**, 2227.
- 7 M. Z. Atashbar and S. Singamaneni, *Sens. Actuators, B*, 2005, **111–112**, 13.
- 8 J. Kong, M. G. Chapline and H. Dai, *Adv. Mater.*, 2001, **13**, 1384.
- 9 Y. Sun and H. H. Wang, *Appl. Phys. Lett.*, 2007, **90**, 213107.
- 10 R. Dasari and F. P. Zamborini, *J. Am. Chem. Soc.*, 2008, **130**, 16138.



**Fig. 7** Capacitance–voltage ( $C$ – $V$ ) curve obtained from our a-IGZO TFT device. For this  $C$ – $V$  measurement, S/D electrodes were used as the ground. Inset band diagrams respectively indicate the depletion and accumulation states of the device.  $V_{\text{FB}}$  and  $V_{\text{Go}}$  were estimated by finding the cross point between two tangent lines, as shown above.

- 11 T. Kiefer, F. Favier, O. Vazquez-Mena, G. Villanueva and J. Brugger, *Nanotechnology*, 2008, **19**, 125502.
- 12 S. Mubeen, B. Yoo and N. V. Myung, *Appl. Phys. Lett.*, 2008, **93**, 133111.
- 13 F. J. Ibanez and F. P. Zamborini, *J. Am. Chem. Soc.*, 2008, **130**, 622.
- 14 E. C. Walter, F. Favier and R. M. Penner, *Anal. Chem.*, 2002, **74**, 1546.
- 15 F. Yang, D. K. Taggart and R. M. Penner, *Nano Lett.*, 2009, **5**, 2177.
- 16 S. Cherevko, N. Kulyk, J. Fu and C. Chung, *Sens. Actuators, B*, 2009, **136**, 338.
- 17 M. Khanuja, S. Kala, B. R. Mehta and F. E. Kruis, *Nanotechnology*, 2009, **20**, 015502.
- 18 T. Kiefer, L. G. Villanueva, F. Fargier, F. Favier and J. Brugger, *Nanotechnology*, 2010, **21**, 505501.
- 19 M. Z. Atashbar, D. Banerji and S. Singamaneni, *IEEE Sens. J.*, 2005, **5**, 792.
- 20 J. Lee, W. Shim, E. Lee, J. Noh and W. Lee, *Angew. Chem., Int. Ed.*, 2011, **50**, 5301.
- 21 J. Lee, J. Noh, S. H. Lee, B. Song, H. Jung, W. Kim and W. Lee, *Int. J. Hydrogen Energy*, 2012, **37**, 7934.
- 22 T. Chang, H. Jung, B. Jang, J. Lee, J. Noh and W. Lee, *Sens. Actuators, A*, 2013, **192**, 140.
- 23 H. Jung, B. Jang, W. Kim, J. Noh and W. Lee, *Sens. Actuators, B*, 2013, **178**, 689.
- 24 R. Feenstra, G. J. Bruin-Hordijk, H. L. M. Bakker, R. Griessen and D. G. Groot, *J. Phys. F: Met. Phys.*, 1983, **13**, L13.
- 25 Y. Sakamoto, K. Takai, I. Takashima and M. Imada, *J. Phys.: Condens. Matter*, 1996, **8**, 3399.
- 26 A. L. Cabrera and R. Aguayo-Soto, *Catal. Lett.*, 1997, **45**, 79.
- 27 Y. Okuhara, Y. Imai, Y. Noguchi and M. Takata, *Bull. Mater. Sci.*, 1999, **22**, 999.
- 28 R. S. Muller, T. I. Kamins and M. Chan, *Device Electronics for Integrated Circuits*, John Wiley & Sons, New York, 3rd edn, 2002, pp. 390–396.